#### **Optical Signal Processing**

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Chapter 5 Imaging and Diffraction-Limited Imaging



#### Introduction

- Imaging is about reproducing the field, or more often the irradiance pattern of an object or scene, at an image plane.
- Geometrical optics, where optical rays are assumed to travel in rectilinear fashion without diffraction, is used extensively in lens and optical system design.
- A proficient approach for image modeling draws on both geometrical optics and diffraction theory.



- In order to form a real image, light from an arbitrary object point must be collected and focused at the image plane.
- For the imaging situation, <u>the lens law describes</u> <u>the relationship needed under the paraxial</u> <u>condition (small ray angles relative to the optical</u> <u>axis) for "best focus" imaging</u>:

$$\frac{1}{z_1} + \frac{1}{z_2} = \frac{1}{f}$$

 f is the lens focal length, z<sub>1</sub> is the distance from the object to the front principal plane of the lens, and z<sub>2</sub> is the distance from the back principal plane to the image location.





- To form a real image, z<sub>1</sub> and z<sub>2</sub> are positive and the lens focal length f is positive.
- A "positive" lens (positive-valued f) converges light rays, whereas a "negative" lens (negative-valued f) diverges rays.
- Practical imaging systems use combinations of lenses to control aberrations, but imaging still requires a positive focal length for the combined lens group.

 The ratio of the image height y<sub>2</sub> to the object height y<sub>1</sub> is known as the *transverse magnification* M<sub>t</sub>, which for a single lens system is given by:

$$M_t = \frac{y_2}{y_1} = -\frac{z_2}{z_1}.$$

- The minus sign indicates an inverted image.
- An imaging system is characterized by its pupils.
- Pupils are virtual apertures that indicate the "opening" available to collect light from the object (entrance pupil) and the "opening" from which the collected light exits on its way to form an image (exit pupil).



- The pupils are images of the physical element in the system, known as the aperture stop, which limits the collection of light.
- The lens is the stop for the system in Fig. 1.
- The stop generates the fundamental diffractive effects in the image: <u>the diffractive effects due to</u> <u>the stop represent the fundamental performance</u> <u>limit of an imaging system.</u>



- Fig. 1-1 illustrates that the physical elements of a system (lenses, mirrors, iris, etc.) can be reduced to entrance pupil (EP) and exit pupil (XP) models.
- The distance from an object point on the optical axis to the EP is  $z_{EP}$  and the distance from the XP to the axial image point is  $z_{XP}$ .
- The entrance pupil diameter is  $D_{EP}$  and the exit pupil diameter is  $D_{XP}$ .



Fig. 1-1. Entrance pupil (EP) and exit pupil (XP) model of an imaging system.

 The important parameter of an imaging system is the *f*-number (*f*/#): the most useful form is the paraxial working *f*/# given by

$$f / \# = \frac{Z_{XP}}{D_{XP}}$$

- <u>The system aperture stop leads to the</u> <u>fundamental diffractive effects in the image.</u>
- For a thin lens imaging system,  $z_1 = z_{EP}$ ,  $z_2 = z_{XP}$ , and  $D_{EP} = D_{XP}$  = lens diameter.



# **Coherent Imaging Theory**



Fig.1-2 shows the general imaging arrangement.



Fig. 1-2. Imaging simulation coordinate definitions.

# **Coherent Imaging Theory**

- Imaging with coherent illumination, such as with a coherent laser, can be described with a convolution operation involving the optical field. It should be noted
- The process is expressed by:

 $U_i(u, v) = h(u, v) \otimes U_{\sigma}(u, v),$ 

- (1)
- that we are only interested in the intensity distribution in the image plane.
- where *u*, *v* are the image plane spatial coordinates, U<sub>i</sub> is the field at the image plane, h is the impulse response for the imaging system, and  $\underline{U}_{a}$  is the ideal geometrical-optics predicted image field, which is a scaled copy of the object field U<sub>o</sub>(x, y):

$$U_g(u, v) = \frac{1}{|M_t|} U_o\left(\frac{u}{M_t}, \frac{v}{M_t}\right).$$

Note that if  $M_t$  is negative, then the resulting image will appear inverted relative to the object.

# **Coherent Imaging Theory**

- In Eq. (1) the ideal geometrical field is "blurred" through the convolution with the impulse function.
- In the frequency domain the corresponding spectra for Eq. (1) are related by:

$$G_i(f_U, f_V) = H(f_U, f_V)G_g(f_U, f_V),$$
(2)

where H is the coherent image transfer function (or amplitude transfer function) and is defined as:

 $H(f_U, f_V) = P(-\lambda z_{XP} f_U, -\lambda z_{XP} f_V), \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_U, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_V, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_V, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_V, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_V, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_V, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_V, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_V, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ [\text{two times FT of } P(\lambda z_{XP} f_V, \lambda z_{XP} f_V)] \quad \text{Refer to Eq. (5-40) in Fourier Optics} \\ \ \text{Refer to Eq. (5-40) in Fourier Optics} \\ \ \text{Refer to Eq. (5-40) in Fourier Optics} \\ \ \text{Refer to Eq. (5-40) in Fourier Optics} \\ \ \text{Refer to Eq. (5-40) in Fourier Optics} \\ \ \text{Refer to Eq. (5-40) in$ 

It should be noted that we are only interested in the

intensity distribution

in the image plane.

- -> inverted]. • where *P* is the pupil function of the system.
- The negative signs in the pupil arguments give a scaled, inverted pupil function.
- In the system with a perfect pupil function, only the boundaries of the pupil are involved with the diffractive effects (diffraction limited).



#### **Coherent Transfer Function Examples**

• For the first example, a square pupil function is:

$$P(x, y) = \operatorname{rect}\left(\frac{x}{2w_{\rm XP}}\right)\operatorname{rect}\left(\frac{y}{2w_{\rm XP}}\right).$$
(3)

• From Eq. (3) the coherent transfer function is:

$$H(f_U, f_V) = \operatorname{rect}\left(\frac{-\lambda z_{XP} f_U}{2w_{XP}}\right) \operatorname{rect}\left(\frac{-\lambda z_{XP} f_V}{2w_{XP}}\right).$$
(4)

- Since the rectangle function is symmetric, the negative signs can be ignored.
- <u>The coherent cutoff frequency along the *u* or *v* (5) <u>direction is defined as</u></u>

$$f_0 = \frac{W_{XP}}{\lambda z_{XP}}$$

• Spatial frequencies with values greater than  $f_0$ will not be preserved in the image plane field.



#### **Coherent Transfer Function Examples**

• A second example is the circular pupil function:

$$P(x, y) = \operatorname{circ}\left(\frac{\sqrt{x^2 + y^2}}{w_{XP}}\right),\tag{6}$$

• The coherent transfer function is

$$H(f_{U}, f_{V}) = \operatorname{circ}\left(\frac{\sqrt{\left(-\lambda z_{XP}f_{U}\right)^{2} + \left(-\lambda z_{XP}f_{V}\right)^{2}}}{w_{XP}}\right), \tag{7}$$
$$H(f_{U}, f_{V}) = \operatorname{circ}\left(\frac{\sqrt{f_{U}^{2} + f_{V}^{2}}}{f_{0}}\right),$$

- where  $f_0$  is again the coherent cutoff frequency as defined in Eq. (5).
- <u>Unlike the square aperture, the cutoff frequency</u> in this case is the same radially in all directions in the frequency plane.



#### **Coherent Transfer Function Examples**

- To observe or record a coherent image, the irradiance given by  $I_i = |U_i|^2$  is actually measured.
- As a result of the squaring operation, the irradiance image can theoretically gain up to *twice* the frequency content of the field—think about the fact that

$$\cos^2\left(2\pi bx\right) = \frac{1}{2} \left[1 - \cos\left(2\pi 2bx\right)\right].$$

• <u>So, when an irradiance image is formed, the</u> following cutoff should be considered:

$$2f_0 = \frac{2w_{XP}}{\lambda z_{XP}}$$

• Given that  $2w_{XP} = D_{XP}$ , and using  $f / \# = \frac{z_{XP}}{D_{XP}}$  $2f_0 = \frac{1}{\lambda(f / \#)}$ 



 One approach for simulating coherent imaging on the computer is based on Eq.(1) and implemented as:

$$U_{i}(u, v) = \Im^{-1} \{ H(f_{U}, f_{V}) \Im \{ U_{g}(u, v) \} \}.$$
(8)

- <u>A simulation begins with a sampled ideal image,</u> <u>usually an image file</u>.
- The sampled ideal image is assumed to have a physical sample interval  $\Delta u$  and side length L.



 <u>The highest spatial frequency available in the</u> <u>ideal image is the Nyquist frequency f<sub>N</sub> = 1/(2Δu),</u> <u>so a diffraction-limited simulation requires</u>

$$2f_0 \le f_N \,. \tag{9}$$

- This condition comes about because you can't model spatial frequencies in the simulation that are not present in the ideal image.
- <u>Applying Eq. (9), substituting for f<sub>N</sub> and rearranging gives a criterion for the sample interval:</u>

$$\Delta u \le \frac{\lambda (f / \#)}{2}. \tag{10}$$



• Given  $L = M\Delta u$ , where M is the number of samples, then  $\lambda(f/\#)$ 

$$L \le M \frac{\pi(j+\pi)}{2}.$$
(11)

- <u>Consider a thin lens with a diameter of 12.5mm,</u> <u>a focal length of 125mm, and in which the</u> <u>wavelength of interest is 0.5µm.</u>
- Assume a 250 × 250 sample array for an imaging simulation.
- First, calculate the *f*-number:  $f/\# = \frac{125}{12.5} = 10$ .
- Inserting the values in Eq. (10) yields

$$\Delta u \le \frac{0.5\mu\mathrm{m}\cdot10}{2} = 2.5\,\mu\mathrm{m}\,,$$

• The side length constraint is:

 $L \le 250 \times 2.5 \,\mu m = 0.625 \,mm$ .

(12)

- The image plane size is limited to 0.625 × 0.625 mm (small area).
- Working with a larger array increases the image size.
- For example, with a 2048 × 2048 sample array, the maximum image plane size would be 5.12×5.12mm—still a relatively small area.



- This illustrates that the Fourier optics-based simulation examines a small part of the image plane for near diffraction-limited performance.
- However, a very large array is needed to model a modest field of view, which might correspond to an image size of, say, 10 or 20 mm in this case.



- Let's work up some code: we use a 250×250 pixel image file that is a USAF 1951 resolution test chart.
- The USAF 1951 is still used today for testing lenses and optical systems.
- Enter the following:

```
coh image Coherent Imaging Example
   용
1
2
3 A=imread('USAF1951B250','png'); %read image file
                      %get image sample size
4 [M,N]=size(A);
                       %reverse row order
5 A=flipud(A);
                   %integer to floating
6 Ig=single(A);
7 Ig=Ig/max(max(Ig)); %normalize ideal image
8
                        %ideal image field
9 uq=sqrt(Iq);
10 L=0.3e-3;
                         %image plane side length (m)
11 du=L/M;
                         %sample interval (m)
12 u=-L/2:du:L/2-du; v=u;
13
14 figure(1)
                         %check ideal image
15 imagesc(u,v,Ig);
16 colormap('gray');
                    xlabel('u (m)'); ylabel('v (m)');
17 axis square
18 axis xy
```



- Since the image file represents an irradiance image, take the square root to get the magnitude of the field (zero phase is assumed in the ideal image).
- The side length is L = 0.3 mm, which satisfies Eq. (12), and with M = 250 the sample interval is Δu = 1.2×10<sup>-6</sup> m.



• The ideal test chart image is shown in Fig. 2.



Fig. 2. USAF test chart ideal image.



 Add the following code to define the imaging system parameters and generate the coherent transfer function:

```
19 lambda=0.5*10^-6; %wavelength
                       %exit pupil radius
20 wxp=6.25e-3;
21 zxp=125e-3; %exit pupil distance
22 f0=wxp/(lambda*zxp); %cutoff frequency
23
24 fu=-1/(2*du):1/L:1/(2*du)-(1/L); %freq coords
25 fv=fu;
26 [Fu,Fv]=meshgrid(fu,fv);
27 H=circ(sqrt(Fu.^2+Fv.^2)/f0);
28
29 figure(2)
                         %check H
30 surf(fu, fv, H.*.99)
31 camlight left; lighting phong
32 colormap('gray')
33 shading interp
34 ylabel('fu (cyc/m)'); xlabel('fv (cyc/m)');
```



- Visible wavelength illumination is assumed.
- For the *f* /10 lens the following are computed:

 $f_N = 1/(2 \cdot 1.2 \times 10^{-6}) = 4.17 \times 10^5$  cycles/m,  $2f_0 = 1/(0.5 \times 10^{-6} \cdot 10) = 2 \times 10^5$  cycles/m,  $f_0 = 1 \times 10^5$  cycles/m.

- Thus, Eq. (9) is satisfied.
- The coherent transfer function is shown in Fig. 3.





- If Eq. (9) were violated, the pupil function in Fig. 3 would reach beyond halfway to the array boundaries.
- Add the following to generate the simulated image:

```
35 H=fftshift(H);
36 Gg=fft2(fftshift(ug));
37 Gi=Gq.*H;
38 ui=ifftshift(ifft2(Gi));
39 Ii=(abs(ui)).^2;
40
41 figure(3)
                          %image result
42 imagesc(u,v,nthroot(Ii,2));
43 colormap('gray'); xlabel('u (m)'); ylabel('v (m)');
44 axis square
45 axis xy
46
47 figure(4)
                        %horizontal image slice
48 vvalue=-0.8e-4; %select row (y value)
49 vindex=round(vvalue/du+(M/2+1)); %convert row index
50 plot(u, Ii(vindex, :), u, Ig(vindex, :), ':');
51 xlabel('u (m)'); ylabel('Irradiance');
```



- <u>Diffraction-limited coherent imaging is</u> <u>implemented and the resulting image is</u> <u>displayed in Fig. 4.</u>
- The features are blurred and some of the threebar groups are unresolved.





- In Fig. 4, the three bars of Group -2, Element 6, appear to be the smallest group that is "resolved."
- Selecting the v-coordinate of  $-0.8 \times 10^{-4}$ m in the profile code gives the display in Fig. 5.
- The large bars of Group –2, Element 1 are clearly resolved but with some obvious ringing effects.
- Element 6 vertical bars are resolved but with much less contrast.





 Figure 6 shows a sequence of image spectra with the coherent transfer function and resulting irradiance images for different pupil sizes (*f*numbers).



Fig. 6. Coherent image spectral magnitude (left column, log scaled) and associated irradiance images (right column).

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